## Introduction to String Theory

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## Exercise Sheet 12

1 In this exercise you will show that the torus partition function of the string

$$A^{0)} \sim \int_{\mathcal{F}} \frac{d^2 \tau}{(\operatorname{Im} \tau)^2} \left( \sqrt{\operatorname{Im} \tau} |\eta(q)|^2 \right)^{-24}, \qquad (1.1)$$

is modular invariant.

(a) Show that the measure

$$\frac{d^2\tau}{(\mathrm{Im}\tau)^2}\,,\tag{1.2}$$

is invariant under fractional linear transformations

$$au o rac{a au + b}{c au + d}$$
. (1.3)

(b) The Dedekind  $\eta$ -function satisfies

$$\eta(\tau+1) = e^{2\pi i/24} \eta(\tau), \qquad \eta(-\frac{1}{\tau}) = \sqrt{-i\tau} \,\eta(\tau).$$
(1.4)

Use this to show that the combination

$$Im\tau |\eta(\tau)|^4 \tag{1.5}$$

is invariant under S and T transformations

$$T: \tau \to \tau + 1, \qquad S: \tau \to -\frac{1}{\tau}.$$
 (1.6)

**2** The action of a string coupled to the background fields G and B is given by

$$S = \frac{1}{4\pi\alpha'} \int_{\Sigma} d^2\sigma \sqrt{g} \left( g^{\alpha\beta} \partial_{\alpha} X^{\mu} \partial_{\beta} X^{\nu} G_{\mu\nu}(X) + i\epsilon^{\alpha\beta} \partial_{\alpha} X^{\mu} \partial_{\beta} X^{\nu} B_{\mu\nu}(X) \right) . \tag{2.1}$$

Show that the action (2.1) is invariant under spacetime diffeomorphisms and B-field gauge transformations

$$B_{\mu\nu} \to B_{\mu\nu} + \partial_{\mu}C_{\nu} - \partial_{\nu}C_{\mu}$$
 (2.2)

**3** (a) Show that the combination

$$H_{\mu\nu\rho} = \partial_{\mu}B_{\nu\rho} + \partial_{\rho}B_{\mu\nu} + \partial_{\nu}B_{\rho\mu}, \qquad (3.1)$$

is invariant under the B-field gauge transformations

$$B_{\mu\nu} \to B_{\mu\nu} + \partial_{\mu}C_{\nu} - \partial_{\nu}C_{\mu}$$
 (3.2)

(b) Let the string worldsheet,  $\Sigma$ , be the boundary of a 3-manifold,  $M_3$ , i.e.  $\Sigma = \partial M_3$ . Show that the string action (2.1) can be written as

$$S = \frac{1}{4\pi\alpha'} \int_{\Sigma} d^2\sigma \sqrt{g} g^{\alpha\beta} \partial_{\alpha} X^{\mu} \partial_{\beta} X^{\nu} G_{\mu\nu}(X) + \frac{i}{12\pi\alpha'} \int_{M_3} d^3\hat{\sigma} \, \eta^{\hat{\alpha}\hat{\beta}\hat{\gamma}} \partial_{\hat{\alpha}} X^{\mu} \partial_{\hat{\beta}} X^{\nu} \partial_{\hat{\gamma}} X^{\rho} H_{\mu\nu\rho}(X) ,$$

$$(3.3)$$

where  $\hat{\sigma}^{\hat{\alpha}}$  are local coordinates on  $M_3$  and  $\eta^{\hat{\alpha}\hat{\beta}\hat{\gamma}}$  is the totally antisymmetric 3-dimensional Levi-Civita tensor density with  $\eta^{123}=1$ .

4 Consider the action

$$S = \frac{1}{2\kappa_0^2} \int d^{26}X \sqrt{-G} e^{-2\Phi} \left( R - \frac{1}{12} H_{\mu\nu\rho} H^{\mu\nu\rho} + 4\partial_{\mu}\Phi \partial^{\mu}\Phi \right). \tag{4.1}$$

Show that its equations of motion for G, B and  $\Phi$  are given by

$$0 = R_{\mu\nu} + 2\nabla_{\mu}\nabla_{\nu}\Phi - \frac{1}{4}H_{\mu\rho\lambda}H_{\nu}^{\rho\lambda},$$

$$0 = -\frac{1}{2}\nabla^{\rho}H_{\mu\nu\rho} + H_{\mu\nu\rho}\nabla^{\rho}\Phi,$$

$$0 = -\frac{1}{2}\nabla^{2}\Phi + \nabla_{\mu}\Phi\nabla^{\mu}\Phi - \frac{1}{24}H_{\mu\nu\rho}H^{\mu\nu\rho}.$$
(4.2)